SECTION 4: MATHEMATICAL MODELING

ESE 330 – Modeling & Analysis of Dynamic Systems

² Introduction

In the last section of notes, we saw how to create a bond graph model from a physical system model.

The next step in the modeling process is the creation of a mathematical model

Mathematical Modeling – Introduction

 You're already familiar with some techniques for creating mathematical models for physical systems



□ First, create a free-body diagram:



Mathematical Modeling – Introduction

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Next, apply Newton's 2nd law



$$\Sigma F = ma$$

$$F_{in}(t) - kx - b\dot{x} = m\ddot{x}$$

rearranging:

$$m\ddot{x} + b\dot{x} + kx = F_{in}(t) \tag{1}$$

This is a *mathematical model*

 A second-order, linear, constant-coefficient, ordinary differential equation

Reduction to a System of 1st-Order ODE's

- Can reduce this 2nd-order ODE to a system of two 1st-order ODE's
- We know that

$$\dot{x} = v$$
 (2)

and

$$\ddot{x} = a = \dot{\nu} \tag{3}$$

Using (2) and (3), rewrite (1), the original ODE

$$m\dot{v} + bv + kx = F_{in}(t)$$
where
$$v = \dot{x}$$
(4)

Reduction to a System of 1st-Order ODE's

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- Equations (4) is a system of first-order ODE's that is equivalent to (1)
- Rearranging (4):

$$\dot{v} = -\frac{k}{m}x - \frac{b}{m}v + \frac{1}{m}F_{in}(t)$$

$$\dot{x} = v$$
(5)

These equations can be put into matrix form :

$$\begin{bmatrix} \dot{x} \\ \dot{\nu} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{k}{m} & -\frac{b}{m} \end{bmatrix} \begin{bmatrix} x \\ \nu \end{bmatrix} + \begin{bmatrix} 0 \\ \frac{1}{m} \end{bmatrix} F_{in}(t)$$
(6)

Reduction to a System of 1st-Order ODE's

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- Let's say we want to consider the *displacement* of the mass as the *output* of the system
- We can add an *output equation* to the mathematical model

$$y = x \tag{7}$$

We can rewrite (7) in a matrix form similar to (6):

$$y = \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} x \\ v \end{bmatrix} + \begin{bmatrix} 0 \end{bmatrix} F_{in}(t)$$
(8)

Mathematical Model

 Together, (6) and (8) comprise the mathematical model for our mechanical system:

$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{k}{m} & -\frac{b}{m} \end{bmatrix} \begin{bmatrix} x \\ v \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \\ \frac{1}{m} \end{bmatrix} F_{in}(t)$$

$$y = \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} x \\ v \end{bmatrix}$$
(9)

Note that x, v, x, v, and y are all *functions of time* The (t) is dropped to simplify the notation
 The convention used here is to only include the (t) for *inputs*, e.g. F_{in}(t)

State-Space Representation

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- The system model of (9) is the state-space representation of the system, or the state-variable equations for the system
- Can be expressed in generic form as

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$$
$$\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$$

(10)

where

- **x**: the state vector
- **• x**: derivative of the state
- **u**: vector of inputs
- **y**: vector of outputs

- A: system matrix
- **B**: input matrix
- **C**: output matrix
- **D**: feed-through matrix

MIMO vs. SISO Systems

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$$

$$\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$$
 (10)

- Note that the state-space model (10) allows for *vectors* of inputs and outputs, u and y
- *Multi-input, multi-output (MIMO) systems* **u** and **y** will be *vectors*
- Single-input, single-output (SISO) systems
 u and *y* will be scalars
- In this course, we'll mostly focus on SISO systems
 For now, we'll assume the more general MIMO case

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System State and State Variables

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- The vector x is the state vector
 - Elements of x are the state variables of the system
- The <u>state</u> of the system is a complete description of the current condition of the system
 - From our energy-based perspective, the state describes all of the energy in a system, i.e. where it is stored, at a given point in time
- The <u>state variables</u> are a (not the) minimum set of system variables required to completely describe the state of a system

State Variables are Not Unique

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- The state vector, i.e. the choice of state variables, for a system is *not unique*
 - In this example, we have chosen displacement and velocity as the state variables, i.e.

$$\mathbf{x} = \begin{bmatrix} x \\ v \end{bmatrix}$$

- Could have chosen other quantities later, we will
- State variables need not even have direct physical significance
- Different state-space representations for the same system are related by *similarity transforms*
 - Beyond the scope of this class

The Feed-Through Matrix

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$$

$$\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$$
 (10)

D is the feed-through or feed-forward matrix Very often zero for physical systems, as in our example

- Non-zero D implies that the input affects the output instantaneously
 - There exists a direct feed-through path from the input to the output

 $\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$ $\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$

- Assume the state space model of (10) represents an nth-order, *m*-input, *p*-output MIMO system
- \Box The *state vector* is an $n \times 1$ column vector
- The system has *m* inputs, so the *input vector* is an $m \times 1$ column vector
- There are p outputs, so the output vector is a $p \times 1$ column vector

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(10)

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$$

$$\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$$
 (10)

 \square If **x** is $n \times 1$, then its derivative, $\dot{\mathbf{x}}$, is also $n \times 1$

The product Ax must have the same dimensions as x, n × 1

\square The system matrix, **A**, is a square $n \times n$ matrix

The product **Bu** must also be n × 1
The vector of inputs, **u**, is m × 1, so **B** is n × m

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$$

$$\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$$
 (10)

 \Box The vector of p outputs, \mathbf{y} , is $p \times 1$

- The product Cx must also have dimension p × 1
 x is n × 1, so C must be p × n
- The product **Du** must also have the same dimension as **y**, *p* × 1
 - **\square** The vector of inputs, **u**, is $m \times 1$, so **D** is $p \times m$

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\square For an *m*-input, *p*-output, MIMO system:

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$$
$$\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$$

Term	Dimension	Те	rm	Dimension
u	$m \times 1$		A	$n \times n$
У	$p \times 1$]	B	$n \times m$
X	$n \times 1$		С	$p \times n$
×	$n \times 1$	1	D	p imes m

L8

 \square For SISO system, u and y, as well as D, are scalars:

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}u$$
$$y = \mathbf{C}\mathbf{x} + Du$$

Term	Dimension	Term	Dimension
u	1×1	Α	$n \times n$
у	1×1	В	$n \times 1$
X	$n \times 1$	С	$1 \times n$
×	$n \times 1$	D	1×1

State-Space Model Explained

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- Remember, our reason for modeling a system is to enable the *analysis of its dynamic behavior*
- Basic idea of the state space model:
 - If the current state of a system is known, and the current and future values of the inputs are known, then the trajectory of the system (i.e. the time-evolution of its state variables) can be determined
 - Don't need explicit knowledge of the history of the system or its inputs no past information
 - All history is accounted for in the current value of the state

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- Consider the physical meaning of the state-space system model

 $\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}\mathbf{u}$ $\mathbf{y} = \mathbf{C}\mathbf{x} + \mathbf{D}\mathbf{u}$

- The time derivative of a system's state variables can be expressed as a linear combination of the current state variables and the current inputs
- The outputs of a system can be expressed as a linear combination of the current state and the current inputs

State-Space Model – Utility

- Again, our goal is to analyze a system's timedomain behavior – the time-evolution of its state variables
- Knowledge of the current state
 variables, as well as the current rate of
 change of those state
 variables, allows us to
 do this



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Where We're Going

- In the previous example, we derived the state-space model for a mechanical system by applying Newton's 2nd law
 - For an electrical system we could have applied Kirchhoff's and Ohm's laws
 - Can always derive a mathematical model by applying domain-specific laws to the physical model
- Our approach will be to *derive state equations from bond-graph system models*

State Equations from Bond-Graph Models

- Bond graphs are *energy-based* models
 - Our choice of state variables will be those that describe the storage of energy within a system at a given instant in time
- State variables will be *energy variables* of the *independent energy-storage elements* in a system
 - Displacements of capacitors
 - Momenta of inertias
- Only independent I's and C's
 - State variables represent a minimum set of system variables needed to completely describe the state

Deriving State Equations from Bond Graphs

Start with the same mechanical system model:



The computational bond graph:



Two independent energy-storage elements

■ State variables will be the energy variables associated with these two elements:

$$\mathbf{x} = \begin{bmatrix} p_2 \\ q_4 \end{bmatrix}$$

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State equation will be of the form:

 $\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}u$ $\begin{bmatrix} \dot{p}_2 \\ \dot{q}_4 \end{bmatrix} = \mathbf{A} \begin{bmatrix} p_2 \\ q_4 \end{bmatrix} + \mathbf{B}e_1(t)$

- In general, state variables will be momenta and displacements
 - Their derivatives will be efforts and flows, respectively
 - For this example:

$$\begin{bmatrix} \dot{p}_2 \\ \dot{q}_4 \end{bmatrix} = \begin{bmatrix} e_2 \\ f_4 \end{bmatrix}$$

State Equation Derivation – Preparation

Annotate the computational bond graph with state variable derivatives

- Efforts on the independent Inertias and the flows on the independent Capacitors
- Apply constitutive laws to annotate the *other power variables* on the *I*'s and *C*'s
- Annotate the *known source power variables* and indicate as functions of time



State Equation Derivation – Procedure

 <u>Objective</u>: derive a set of n equations, each expressing a state variable derivative as a linear combination of state variables and inputs

Determine the A and B matrices

First, choose a state variable and write its derivative as an effort or flow:

$$\dot{p}_2 = e_2 \tag{1}$$

- Next, use the causality assigned to the bond graph to work from (1) to a state equation
 - Express \dot{p}_2 as a linear combination of states and inputs
 - Will ultimately relate an effort or flow to a state variable by applying a constitutive relationship for an energy-storage element

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e₂ is an effort on a 1-jct
 ■ Caused by e₁, e₃, and e₄, and e₁ is known, so

$$\dot{p}_2 = e_2 = e_1(t) - e_3 - e_4$$
 (2)

Relate e₃ to f₃ using the const. law for the resistor

$$e_3 = R_3 f_3 \tag{3}$$

□ f_3 is the flow on a 1-jct, set by f_2 , related to s.v. p_2 by the const. law for the inertia

$$f_3 = f_2 = \frac{1}{I_2} p_2 \tag{4}$$

R

e₁(t)

State Equation Derivation – Procedure

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Substituting (4) into (3)

$$e_3 = \frac{R_3}{I_2} p_2 \tag{5}$$

And substituting (5) back into (2)

$$\dot{p}_2 = e_1(t) - \frac{R_3}{I_2} p_2 - e_4 \tag{6}$$

(7)

- □ Still need to eliminate e_4
 - *e*₄ related to state variable *q*₄ through constitutive law for the capacitor

$$e_4 = \frac{1}{C_4}q_4$$



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- Substituting (7) into (6) yields the first of two state equations

$$\dot{p}_2 = -\frac{R_3}{I_2} p_2 - \frac{1}{C_4} q_4 + e_1(t)$$

 \Box Next, follow a similar procedure for q_4

$$\dot{q}_4 = f_4 \tag{9}$$

□ f_4 is the flow on a 1-jct, set by f_2 , related to state variable p_2 by the const. law for the inertia

$$f_4 = f_2 = \frac{1}{I_2} p_2 \tag{10}$$

(8)

К

3

e1(t)

Substituting (10) into (9) yields the second of two state equations

$$\dot{q}_4 = \frac{1}{I_2} p_2 \tag{11}$$

Combine (8) and (11) into the *state-variable model* for our system in matrix form

$$\begin{bmatrix} \dot{p}_2 \\ \dot{q}_4 \end{bmatrix} = \begin{bmatrix} -\frac{R_3}{I_2} & -\frac{1}{C_4} \\ \frac{1}{I_2} & 0 \end{bmatrix} \begin{bmatrix} p_2 \\ q_4 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} e_1(t)$$
(12)



 Can now replace the computational bond graph parameters in (12) with physical system parameters

$$\begin{bmatrix} \dot{p} \\ \dot{\chi} \end{bmatrix} = \begin{bmatrix} -\frac{b}{m} & -k \\ \frac{1}{m} & 0 \end{bmatrix} \begin{bmatrix} p \\ \chi \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} F_{in}(t)$$
(13)

State Equation Derivation – Output Equation

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- Can also define an *output equation* as part of our state-space model
- Suppose we want to consider the velocity of the mass as our output

- $k \longrightarrow V$ $\downarrow \qquad b \qquad m \rightarrow F_{in}(t)$
- Constitutive relation relates an inertia's flow to its momentum:

$$f_2 = \nu = \frac{1}{I_2} p_2 = \frac{1}{m} p \tag{14}$$

The output equation would be:

$$y = \begin{bmatrix} 1/m & 0 \end{bmatrix} \begin{bmatrix} p \\ x \end{bmatrix}$$
(15)

Equations (13) and (15) comprise the complete state-space system model

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State Equation Derivation – Output Equation

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- Perhaps, instead, we want to consider the *displacement of the mass* as our output
 - Same as spring displacement a state variable

 $y = \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{vmatrix} p \\ \gamma \end{vmatrix}$

- $k \longrightarrow K$ $\downarrow \downarrow \downarrow b \qquad m \rightarrow F_{in}(t)$
- State-space model, including output equation, becomes:

$$\begin{bmatrix} \dot{p} \\ \dot{\chi} \end{bmatrix} = \begin{bmatrix} -\frac{b}{m} & -k \\ \frac{1}{m} & 0 \end{bmatrix} \begin{bmatrix} p \\ \chi \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix} F_{in}(t)$$

(16)

State Equation Derivation – Causality

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- In this example, assignment of causality yielded the simplest result:
 - All energy-storage elements ended up in integral causality

 all were independent
 - No resistors had their causality arbitrarily assigned
- Lack of derivative causality and/or algebraic loops (resistor fields) results in straightforward state equation derivation
 - Unfortunately, the inverse is also true
- Next, we'll look at two more examples without derivative causality or algebraic loops
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- Consider the mechanical example from Section 3
- Four independent energy-storage elements
 - Fourth-order system
 - Four state variables:

$$\mathbf{x} = \begin{bmatrix} p_3 \\ q_5 \\ q_8 \\ p_{10} \end{bmatrix}$$



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State Equation Derivation – Example 1

- Annotate the bond graph:
 - State variable derivatives
 - Efforts on independent inertias
 - Flows on independent capacitors
 - Use constitutive laws and state variables to express:
 - Flows on independent inertias
 - Efforts on independent capacitors

Known source quantities



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- Choose a state variable derivative and express it as an effort or a flow

$$\dot{p}_3 = e_3 = e_1(t) - e_2(t) - e_4$$
 (1)

Known source efforts can remain
 Need to eliminate e₄
 Effort on a 0-jct, set by e₅

$$e_4 = e_5 = \frac{1}{C_5}q_5$$

Substituting (2) into (1) yields the *first* of four state equations

$$\dot{p}_3 = -\frac{1}{c_5}q_5 + e_1(t) - e_2(t)$$



(2)

(3)

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Move on to the next state variable

$$\dot{q}_5 = f_5 = f_4 - f_6 \tag{4}$$

□ f_4 and f_6 are both flows on 1-jct's set by f_3 and f_{10} , respectively

$$f_4 = f_3 = \frac{1}{I_3} p_3 \tag{5}$$

$$f_6 = f_{10} = \frac{1}{I_{10}} p_{10} \tag{6}$$

Substituting (6) and (5) into (4) yields the *second state equation*

$$\dot{q}_5 = \frac{1}{I_3} p_3 - \frac{1}{I_{10}} p_{10} \tag{7}$$



(8)

(9)

 \Box Move on to \dot{q}_8

 $\dot{q}_8 = \frac{1}{I_{10}} p_{10}$

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$$\dot{q}_8 = f_8 = f_{10} = \frac{1}{I_{10}} p_{10}$$

which gives the *third state equation*

 \Box Finally, derive the equation for \dot{p}_{10}

$$\dot{p}_{10} = e_{10} = e_6 - e_7(t) - e_8 - e_9$$
 (10)

 \Box e_6 is the effort on a 0-jct, set by e_5

$$e_6 = e_5 = \frac{1}{C_5} q_5 \tag{11}$$

$$\begin{array}{c}
C & \frac{1}{C_{8}} \cdot q_{8} & g \\
S_{e}^{(t)} & 1 & \dot{p}_{10} \\
S_{e}^{(t)} & 1 & \frac{1}{I_{10}} \cdot p_{10} \\
f & \frac{1}{I_{10}} \cdot p_{10} \\
f & \frac{1}{I_{10}} \cdot q_{5} \\
f & 0 & \frac{1}{C_{5}} \cdot$$

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State Equation Derivation – Example 1

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- \Box e_8 is related to state variable q_8
 - $e_8 = \frac{1}{C_8} q_8 \tag{12}$
- \Box e_9 can be related to f_9 using the constitutive law for resistor R_9

$$e_9 = R_9 f_9 \tag{13}$$

□ And, f_9 is the flow on a 1-jct, set by f_{10}

$$e_9 = R_9 f_{10} = R_9 \frac{1}{I_{10}} p_{10} \tag{14}$$

Substituting (11), (12), and (14) into (10) yields the *final state equation*

$$\dot{p}_{10} = \frac{1}{C_5} q_5 - \frac{1}{C_8} q_8 - \frac{R_9}{I_{10}} p_{10} - e_7(t) \tag{15}$$

$$\begin{array}{c}
C & \frac{1}{C_{8}} \cdot q_{8} & g \\
S_{e} & \frac{1}{C_{8}} \cdot q_{8} & g \\
S_{e} & \frac{1}{C_{8}} \cdot q_{8} & g \\
S_{e} & \frac{1}{C_{5}} \cdot q_{1} & f_{10} & f_{10} \\
& & 1 & \frac{1}{I_{10}} \cdot p_{10} & f_{10} \\
& & 0 & \frac{1}{C_{5}} \cdot q_{5} & f_{10} \\
& & 0 & \frac{1}{C_{5}} \cdot q_{5} & f_{10} \\
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& & 0 & \frac{1}{C_{5}} \cdot q_{10} & f_{10} & f_{10} & f_{10} \\
& & 0 & \frac{1}{C_{5}} \cdot q_{10} & f_{10} & f_{10} & f_{10} & f_{10} & f_{10} & f_{10} \\
& & 0 & \frac{1}{C_{5}} \cdot q_{10} & f_{10} &$$

Combine the state equations into matrix form

$$\begin{bmatrix} \dot{p}_{3} \\ \dot{q}_{5} \\ \dot{q}_{8} \\ \dot{p}_{10} \end{bmatrix} = \begin{bmatrix} 0 & -\frac{1}{c_{5}} & 0 & 0 \\ \frac{1}{I_{3}} & 0 & 0 & -\frac{1}{I_{10}} \\ 0 & 0 & 0 & \frac{1}{I_{10}} \\ 0 & \frac{1}{C_{5}} & -\frac{1}{C_{8}} & -\frac{R_{9}}{I_{10}} \end{bmatrix} \begin{bmatrix} p_{3} \\ q_{5} \\ q_{8} \\ p_{10} \end{bmatrix} + \begin{bmatrix} 1 & -1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} e_{1}(t) \\ e_{2}(t) \\ e_{7}(t) \end{bmatrix}$$
(16)

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State Equation Derivation – Example 1

- Let the *position of each mass* to be our outputs
 Two outputs
- Displacement of m₂ (I₁₀)
 is the displacement of the upper spring

$$x_2 = q_8 \tag{17}$$

 Displacement of m₁ is the sum of the spring displacements

$$x_1 = q_5 + q_8$$
 (18)



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- Combine (17) and (18) into our output equation
 Multiple outputs, so C will be a *matrix*
- Complete state-space model, including output equation:

$$\dot{\mathbf{x}} = \begin{bmatrix} \dot{p}_3 \\ \dot{q}_5 \\ \dot{q}_8 \\ \dot{p}_{10} \end{bmatrix} = \begin{bmatrix} 0 & -\frac{1}{c_5} & 0 & 0 \\ \frac{1}{l_3} & 0 & 0 & -\frac{1}{l_{10}} \\ 0 & 0 & 0 & \frac{1}{l_{10}} \\ 0 & \frac{1}{c_5} & -\frac{1}{c_8} & -\frac{R_9}{l_{10}} \end{bmatrix} \begin{bmatrix} p_3 \\ q_5 \\ q_8 \\ p_{10} \end{bmatrix} + \begin{bmatrix} 1 & -1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} e_1(t) \\ e_2(t) \\ e_7(t) \end{bmatrix}$$

$$\begin{bmatrix} \dot{p}_{10} \end{bmatrix} \begin{bmatrix} 0 & \frac{1}{c_5} & -\frac{1}{c_8} & -\frac{R_9}{I_{10}} \end{bmatrix} \begin{bmatrix} p_{10} \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 0 & \frac{1}{c_5} & -\frac{1}{c_8} & -\frac{R_9}{I_{10}} \end{bmatrix}$$
$$\mathbf{y} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 1 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} p_3 \\ q_5 \\ q_8 \\ p_{10} \end{bmatrix}$$



(21)

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- Can rewrite our state-space model, substituting in physical parameters
 - q₁ and q₂ are the displacements of springs k₁ and k₂, respectively

$$\dot{\mathbf{x}} = \begin{bmatrix} \dot{p}_1 \\ \dot{q}_1 \\ \dot{q}_2 \\ \dot{p}_2 \end{bmatrix} = \begin{bmatrix} 0 & -k_1 & 0 & 0 \\ \frac{1}{m_1} & 0 & 0 & -\frac{1}{m_2} \\ 0 & 0 & 0 & \frac{1}{m_2} \\ 0 & k_1 & -k_2 & -\frac{b}{m_2} \end{bmatrix} \begin{bmatrix} p_1 \\ q_1 \\ q_2 \\ p_2 \end{bmatrix} + \begin{bmatrix} 1 & -1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} F_{in}(t) \\ m_1g \\ m_2g \end{bmatrix}$$

$$\mathbf{y} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 1 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} p_1 \\ q_1 \\ q_2 \\ p_2 \end{bmatrix}$$

(21)



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- A slightly modified version of the electrical circuit from Section 3:



The computational bond graph for this circuit:



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- Three independent energy-storage elements
 Third order
- State variables:

$$\mathbf{x} = \begin{bmatrix} p_3 \\ q_5 \\ q_9 \end{bmatrix}$$

 Annotate the computational bond graph





 $\,\,$ Begin with equation for \dot{p}_3

$$\dot{p}_3 = e_3 = e_1(t) - e_2 - e_4$$

$$e_2 = R_2 f_2 = R_2 f_3 = R_2 \frac{1}{I_3} p_3 \qquad (2$$



 $\Box e_4$ is the effort on a 0-jct, set by the effort on C_5

$$e_4 = e_5 = \frac{1}{C_5} q_5 \tag{3}$$

 Substituting (2) and (3) into (1) gives the *first of three state* equations

$$\dot{p}_3 = -\frac{R_2}{I_3} p_3 - \frac{1}{C_5} q_5 + e_1(t) \tag{4}$$

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□ The transformer modulus relates f_6 to f_7 , which is the flow on a 1-jct, set by f_8

$$f_6 = \frac{N_2}{N_1} f_7 = \frac{N_2}{N_1} f_8 = \frac{N_2}{N_1} \frac{1}{R_8} e_8$$
(7)

 $\Box e_8$ is algebraically related to e_7 and e_9

$$e_8 = e_7 - e_9 = e_7 - \frac{1}{C_9}q_9 \tag{8}$$

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The transformer relates e_7 back to e_6 , which is set by e_5

$$e_7 = \frac{N_2}{N_1} e_6 = \frac{N_2}{N_1} e_5 = \frac{N_2}{N_1} \frac{1}{C_5} q_5$$
 (9)

Substituting (9) into (8) gives

$$e_8 = \frac{N_2}{N_1} \frac{1}{C_5} q_5 - \frac{1}{C_9} q_9$$



$$f_6 = \frac{N_2}{N_1} \frac{1}{R_8} \left(\frac{N_2}{N_1} \frac{1}{C_5} q_5 - \frac{1}{C_9} q_9 \right)$$
(11)

Using (11) and (6) in (5) gives us our *second state equation*

$$\dot{q}_5 = \frac{1}{I_3} p_3 - \left(\frac{N_2}{N_1}\right)^2 \frac{1}{R_8 C_5} q_5 + \frac{N_2}{N_1} \frac{1}{R_8 C_9} q_9 \tag{12}$$

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(10)



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 \Box Finally, derive the equation for \dot{q}_9

$$\dot{q}_9 = f_9 \tag{13}$$

□ f_9 is the flow on a 1-jct, which is set by f_8

$$f_9 = f_8 = \frac{1}{R_8}e_8$$



Substituting (10) into (14)

$$f_9 = \frac{1}{R_8} \left(\frac{N_2}{N_1} \frac{1}{C_5} q_5 - \frac{1}{C_9} q_9 \right)$$
(15)

Substituting (15) in (13) gives us our *third state equation*

$$\dot{q}_9 = \frac{N_2}{N_1} \frac{1}{R_8 C_5} q_5 - \frac{1}{R_8 C_9} q_9 \tag{16}$$

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Combine the state equations in matrix form

$$\dot{\mathbf{x}} = \begin{bmatrix} \dot{p}_3 \\ \dot{q}_5 \\ \dot{q}_9 \end{bmatrix} = \begin{bmatrix} -\frac{R_2}{I_3} & -\frac{1}{C_5} & 0 \\ \frac{1}{I_3} & -\left(\frac{N_2}{N_1}\right)^2 \frac{1}{R_8 C_5} & \frac{N_2}{N_1} \frac{1}{R_8 C_9} \\ 0 & \frac{N_2}{N_1} \frac{1}{R_8 C_5} & -\frac{1}{R_8 C_9} \end{bmatrix} \begin{bmatrix} p_3 \\ q_5 \\ q_9 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} e_1(t)$$
(17)

 Replacing computational bond graph parameters with physical parameters

$$\dot{\mathbf{x}} = \begin{bmatrix} \dot{\lambda}_1 \\ \dot{q}_1 \\ \dot{q}_2 \end{bmatrix} = \begin{bmatrix} -\frac{R_1}{L_1} & -\frac{1}{C_1} & 0 \\ \frac{1}{L_1} & -\left(\frac{N_2}{N_1}\right)^2 \frac{1}{R_3 C_1} & \frac{N_2}{N_1} \frac{1}{R_3 C_2} \\ 0 & \frac{N_2}{N_1} \frac{1}{R_3 C_1} & -\frac{1}{R_3 C_2} \end{bmatrix} \begin{bmatrix} \lambda_1 \\ q_1 \\ q_2 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} v_{in}(t)$$
(18)



Choosing the voltage across C_2 as our output, the complete statespace system representation is

$$\dot{\mathbf{x}} = \begin{bmatrix} \dot{\lambda}_1 \\ \dot{q}_1 \\ \dot{q}_2 \end{bmatrix} = \begin{bmatrix} -\frac{R_1}{L_1} & -\frac{1}{C_1} & 0 \\ \frac{1}{L_1} & -\left(\frac{N_2}{N_1}\right)^2 \frac{1}{R_3 C_1} & \frac{N_2}{N_1 R_3 C_2} \\ 0 & \frac{N_2}{N_1 R_3 C_1} & -\frac{1}{R_3 C_2} \end{bmatrix} \begin{bmatrix} \lambda_1 \\ q_1 \\ q_2 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} v_{in}(t)$$

$$y = v_d = \begin{bmatrix} 0 & 0 & 1/C_2 \end{bmatrix} \begin{bmatrix} \lambda_1 \\ q_1 \\ q_2 \end{bmatrix}$$
(19)



- Instead let the voltage across L₁ be the system output
 That is, the effort associated with L₁
 - Effort is the time derivative of momentum, so

$$y = v_{L1} = v_a - v_b = \dot{\lambda}_1$$
 (20)

□ The output equation can be extracted from (19)

$$y = v_{L1} = \begin{bmatrix} -\frac{R_1}{L_1} & -\frac{1}{C_1} & 0 \end{bmatrix} \begin{bmatrix} \lambda_1 \\ q_1 \\ q_2 \end{bmatrix} + v_{in}(t)$$
(21)

□ Note that, in this case, the feed-through term, *D*, is non-zero



Consider the following electrical circuit



- Causality assignment is completed by arbitrarily assigning the causality of resistor R_2 (or R_6)
 - System contains an *algebraic loop* (resistor field)

$$S_{e} \xrightarrow{e_{1}(t)} 1 \xrightarrow{3} 0 \xrightarrow{5} 1 \xrightarrow{\frac{1}{C_{7}} \cdot q_{7}} C$$

$$\stackrel{2}{\downarrow} \stackrel{\dot{p}_{4}}{\downarrow} \stackrel{\frac{1}{I_{4}} \cdot p_{4}}{I} \xrightarrow{6} \xrightarrow{R} R$$

 Presence of the algebraic loop will complicate the state equation derivation a bit

Second-order systemState variables are:

$$\mathbf{x} = \begin{bmatrix} p_4 \\ q_7 \end{bmatrix}$$



Begin deriving equations as usual

$$\dot{p}_4 = e_4 = e_3 = e_1(t) - e_2 = e_1(t) - R_2 f_2$$
 (1)

$$f_2 = f_3 = \frac{1}{I_4}p_4 + f_5 = \frac{1}{I_4}p_4 + f_6$$
(2)

$$f_6 = \frac{1}{R_6} e_6 = \frac{1}{R_6} \left(e_5 - \frac{1}{C_7} q_7 \right)$$
(3)

$$f_6 = \frac{1}{R_6} \left(e_3 - \frac{1}{C_7} q_7 \right) \tag{4}$$

e₃ has reentered the formulation, and we're back where we started in (1)
 An *algebraic loop*

Algebraic Loops – Procedure

- 59
- 1. The **output** of the resistor whose causality was arbitrarily assigned $-e_2$ in this case, though f_6 would work equally well is the **auxiliary variable**
- 2. Derive an expression relating the *auxiliary variable* to the *state variables, inputs,* and to *itself*
- 3. Proceed with the *state equation derivation* as usual, but leave the auxiliary variable in the formulation along with state variables and inputs
- 4. *Substitute* the result from step 2 into the result from step 3
- One auxiliary variable for each algebraic loop present
 Multiple loops require solution of a system of equations
- Apply this procedure first, whenever causality assignment involves an arbitrary assignment of resistor causality

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- Follow causality to derive an expression for auxiliary variable e₂

$$e_2 = R_2 f_2 = R_2 f_3 = R_2 \left(\frac{1}{I_4} p_4 + f_5\right)$$
 (5)

$$f_5 = f_6 = \frac{1}{R_6} e_6 = \frac{1}{R_6} \left(e_5 - \frac{1}{C_7} q_7 \right)$$
(6)

$$e_5 = e_3 = e_1(t) - e_2 \tag{7}$$

S_e-

e₂ is the aux. variable, so it can remain in the expression
 Substituting (7) into (6) into (5)

$$e_2 = \frac{R_2}{I_4} p_4 + \frac{R_2}{R_6} e_1(t) - \frac{R_2}{R_6} e_2 - \frac{R_2}{R_6 C_7} q_7$$
(8)

$$e_2 \frac{R_2 + R_6}{R_6} = \frac{R_2}{I_4} p_4 - \frac{R_2}{R_6 C_7} q_7 + \frac{R_2}{R_6} e_1(t)$$
(9)

 $\dot{p}_4 \left| \frac{1}{l_4} \cdot p_4 \right|$

6

2

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Solve (9) for e_2

$$e_2 = \frac{R_2 R_6}{R_2 + R_6} \frac{1}{I_4} p_4 - \frac{R_2}{(R_2 + R_6)C_7} q_7 + \frac{R_2}{R_2 + R_6} e_1(t)$$
(10)

 Now, whenever e₂ appears in the formulation, substitute in the expression in (10)



□ Going back to (1), we had

$$\dot{p}_4 = e_1(t) - e_2 \tag{1}$$

Substituting in (10) yields the *first state equation*

$$\dot{p}_4 = -\frac{R_2 R_6}{R_2 + R_6} \frac{1}{I_4} p_4 + \frac{R_2}{(R_2 + R_6)C_7} q_7 + \frac{R_6}{R_2 + R_6} e_1(t)$$
(11)

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Moving on to \dot{q}_7

$$\dot{q}_7 = f_7 = f_6 \tag{12}$$

 $S_{e} \xrightarrow{e_{1}(t)} 1 \xrightarrow{3} 0 \xrightarrow{5} 1 \xrightarrow{\frac{1}{C_{7}} \cdot q_{7}} C$ $\stackrel{2}{\downarrow} \stackrel{\dot{p}_{4}}{\downarrow} \frac{1}{l_{4}} \cdot p_{4} \xrightarrow{6} \qquad R \qquad I \qquad R$

• We already have an expression for f_6 in (6) and (7)

$$\dot{q}_7 = \frac{1}{R_6} e_1(t) - \frac{1}{R_6} e_2 - \frac{1}{R_6 C_7} q_7 \tag{13}$$

 \Box Substituting in (10) to eliminate e_2

$$\dot{q}_7 = \frac{1}{R_6} e_1(t) - \frac{R_2}{R_2 + R_6} \frac{1}{I_4} p_4 + \frac{R_2}{(R_2 + R_6)R_6C_7} q_7 - \frac{R_2}{(R_2 + R_6)R_6} e_1(t) - \frac{1}{R_6C_7} q_7$$

Rearranging gives the *second state equation*

$$\dot{q}_7 = -\frac{R_2}{R_2 + R_6} \frac{1}{I_4} p_4 - \frac{1}{(R_2 + R_6)C_7} q_7 + \frac{1}{R_2 + R_6} e_1(t)$$
(14)



 Assembling (11) and (14) in matrix form gives our state variable system model

$$\begin{bmatrix} \dot{p}_4\\ \dot{q}_7 \end{bmatrix} = \begin{bmatrix} -\frac{R_2 R_6}{R_2 + R_6} \frac{1}{I_4} & \frac{R_2}{(R_2 + R_6)C_7} \\ -\frac{R_2}{R_2 + R_6} \frac{1}{I_4} & -\frac{1}{(R_2 + R_6)C_7} \end{bmatrix} \begin{bmatrix} p_4\\ q_7 \end{bmatrix} + \begin{bmatrix} \frac{R_6}{R_2 + R_6} \\ \frac{1}{R_2 + R_6} \end{bmatrix} e_1(t)$$
(15)

 $\lfloor q \rfloor$



 Substitute in physical parameters and define an output equation for the voltage across the capacitor, v_b

$$\begin{bmatrix} \dot{\lambda} \\ \dot{q} \end{bmatrix} = \begin{bmatrix} -\frac{R_1 R_2}{R_1 + R_2} \frac{1}{L} & \frac{R_1}{(R_1 + R_2)C} \\ -\frac{R_1}{R_1 + R_2} \frac{1}{L} & -\frac{1}{(R_1 + R_2)C} \end{bmatrix} \begin{bmatrix} \lambda \\ q \end{bmatrix} + \begin{bmatrix} \frac{R_2}{R_1 + R_2} \\ \frac{1}{R_1 + R_2} \end{bmatrix} v_{in}(t)$$
$$y = \begin{bmatrix} 0 & 1/C \end{bmatrix} \begin{bmatrix} \lambda \\ q \end{bmatrix}$$

(16)

Next, consider a mechanical system



Causality assignment is completed by arbitrarily assigning the causality of resistor R_2 (or R_4)



 A very similar bond graph to the electrical circuit in the previous example

- 66
- A *second-order system* with state variables:

$$\mathbf{x} = \begin{bmatrix} q_1 \\ p_6 \end{bmatrix}$$

 A second-order system with state variables:

$$\mathbf{x} = \begin{bmatrix} q_1 \\ p_6 \end{bmatrix}$$

- An *algebraic loop* is present, so we'll immediately go to the procedure outlined in the previous example
- \Box **Auxiliary variable** is f_2
 - Express f_2 in terms of state variables, inputs, and itself

$$f_2 = \frac{1}{R_2} e_2 = \frac{1}{R_2} (e_3 - e_1) = \frac{1}{R_2} \left(e_4 - \frac{1}{C_1} q_1 \right) \tag{1}$$

$$e_4 = R_4 f_4 = R_4 (f_5 - f_3) = R_4 \left(\frac{1}{I_6} p_6 - f_2\right)$$
(2)

 \Box f_2 is the auxiliary variable, so it can remain in the expression

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Substitute (2) into (1)

$$f_2 = \frac{R_4}{R_2} \frac{1}{I_6} p_6 - \frac{R_4}{R_2} f_2 - \frac{1}{R_2 C_1} q_1$$

□ Then solve for f_2

$$f_2\left(\frac{R_2+R_4}{R_2}\right) = -\frac{1}{R_2C_1}q_1 + \frac{R_4}{R_2}\frac{1}{I_6}p_6$$
(4)

$$f_2 = -\frac{1}{(R_2+R_4)C_1}q_1 + \frac{R_4}{R_2+R_4}\frac{1}{I_6}p_6$$
(5)

(3) $C \xrightarrow{\frac{1}{C_{1}} \cdot q_{1}}{\dot{q}_{1}} 1 \xrightarrow{\gamma_{3}} 0 \xrightarrow{5} 1 \xrightarrow{\gamma_{6}}{1} \xrightarrow{\rho_{6}}{1}$

Now, Proceed with the state equation derivation

• Whenever the auxiliary variable, f_2 , appears in the formulation, it will be replaced with (5)

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 $\,\,$ Start with \dot{q}_1

$$\dot{q}_1 = f_1 = f^2$$
 (6)

 Substituting (5) into (6) gives the *first state equation*

$$\dot{q}_1 = -\frac{1}{(R_2 + R_4)C_1}q_1 + \frac{R_4}{R_2 + R_4}\frac{1}{I_6}p_6$$

 \square Moving on to \dot{p}_6

$$\dot{p}_6 = e_6 = e_7(t) - e_5 = e_7(t) - e_4$$
(8)

$$e_4 = R_4 f_4 = R_4 (f_5 - f_3) = R_4 \left(\frac{1}{I_6} p_6 - f_2\right)$$
(9)

(7)

 $C \xrightarrow{\frac{1}{C_{1}} \cdot q_{1}} 1 \xrightarrow{4} 0 \xrightarrow{5} 1 \xrightarrow{1} S_{e}$

Substitute (9) into (8)

$$\dot{p}_6 = e_7(t) - \frac{R_4}{I_6}p_6 + R_4f_2$$
 (2)

Substituting (5) in for f₂ gives the second state equation



$$\dot{p}_{6} = e_{7}(t) - \frac{R_{4}}{I_{6}}p_{6} + R_{4}\left(-\frac{1}{(R_{2}+R_{4})C_{1}}q_{1} + \frac{R_{4}}{R_{2}+R_{4}}\frac{1}{I_{6}}p_{6}\right)$$
(11)
$$\dot{p}_{6} = -\frac{R_{4}}{(R_{2}+R_{4})C_{1}}q_{1} - \frac{R_{2}R_{4}}{R_{2}+R_{4}}\frac{1}{I_{6}}p_{6} + e_{7}(t)$$
(12)

□ In matrix form:

$$\begin{bmatrix} \dot{q}_1 \\ \dot{p}_6 \end{bmatrix} = \begin{bmatrix} -\frac{1}{(R_2 + R_4)C_1} & \frac{R_4}{R_2 + R_4} \frac{1}{I_6} \\ -\frac{R_4}{(R_2 + R_4)C_1} & -\frac{R_2R_4}{R_2 + R_4} \frac{1}{I_6} \end{bmatrix} \begin{bmatrix} q_1 \\ p_6 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} e_7(t)$$
(13)

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- Note that the origin of the algebraic loop in this example was a *modeling* assumption
 - The connection point between the spring and dampers was considered *massless*
 - Instead we could account for the mass of this junction



Now, there are no arbitrary causality assignments and *no algebraic loops*



State equation derivation will be greatly simplified

- 71
- System is now *third-order*, due to the additional independent energy-storage element



State equation, after replacing physical parameters:

$$\begin{bmatrix} \dot{x}_2 \\ \dot{p}_1 \\ \dot{p}_2 \end{bmatrix} = \begin{bmatrix} 0 & 0 & \frac{1}{m_2} \\ 0 & -\frac{b_2}{m_1} & \frac{b_2}{m_1} \\ -k & \frac{b_2}{m_1} & -\frac{b_1+b_2}{m_2} \end{bmatrix} \begin{bmatrix} x_2 \\ p_1 \\ p_2 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} F_{in}(t)$$
(14)

Looks very different from the original second order model, but if $m_2 \ll m_1$, their behaviors are nearly identical

72 Derivative Causality
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 Consider the mechanical system from Section 3

The computational bond graph:



- Two *independent* energy-storage elements
 - Second-order
 - State variables are:

$$\mathbf{x} = \begin{bmatrix} p_2 \\ q_7 \end{bmatrix}$$



- I₆ is in *derivative Causality*
 - Not independent

Does not contribute a state

Its energy variable p₆ (would be a q for an C-element) is algebraically related to the state variables



 Annotate the bond graph
 Include the energy variable annotation for the dependent inertia





 $\square p_6$ is not a state variable

- State equation derivation requires first determining the algebraic relationship between p₆ and the state variables, p₂ and q₇
- When p_6 or \dot{p}_6 enters the formulation, substitute in this relationship or its derivative

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Derivative Causality – Procedure

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- For the dependent energy-storage element, *apply the constitutive law 'backwards'* - i.e. express the energy variable as a function of a power variable
 - a. <u>Inertia</u>: express momentum as a function of flow
 - b. <u>Capacitor</u>: express displacement as a function of effort
- 2. Follow causality to relate that power variable to the state variables and inputs
- 3. Substitute the expression from step 2 into that from step 1
- 4. When the energy variable (or its derivative) enters the formulation, *substitute* in the expression from step 3

- 77
- Apply the constitutive law for *I*₆ 'backwards'
 Express *n* as a function of *f*
 - Express p_6 as a function of f_6

$$p_6 = I_6 f_6$$
 (1)



 \Box Follow causality to express f_6 in terms of state variables and inputs

$$f_6 = f_4 = \frac{b}{a} f_3 = \frac{b}{a} \frac{1}{I_2} p_2 \tag{2}$$

Substituting (2) into (1)

$$p_6 = \frac{b}{a} \frac{I_6}{I_2} p_2 \tag{3}$$

Now proceed with derivation, using (3) when needed

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- Begin state equation derivation with \dot{p}_2 $\dot{p}_2 = e_2 = e_1(t) - e_3$ (4) TF relates e_3 to e_4 $\dot{p}_2 = e_1(t) - \frac{b}{a}e_4 = e_1(t) - \frac{b}{a}\left(e_5 + \dot{p}_6 + \frac{1}{C_7}q_7\right)$ (5) $e_5 = R_5f_5 = R_5f_4 = R_5\frac{b}{a}f_3 = R_5\frac{b}{a}\frac{1}{I_2}p_2$ (6)
- Substituting (6) into (5)

$$\dot{p}_{2} = e_{1}(t) - \frac{b}{a} \left(\frac{b}{a} \frac{R_{5}}{I_{2}} p_{2} + \dot{p}_{6} + \frac{1}{C_{7}} q_{7} \right)$$

$$\dot{p}_{2} = e_{1}(t) - \left(\frac{b}{a} \right)^{2} \frac{R_{5}}{I_{2}} p_{2} - \frac{b}{a} \dot{p}_{6} - \frac{b}{a} \frac{1}{C_{7}} q_{7}$$
(7)

A *p*₆ term has entered the formulation

D Differentiate (3)

$$\dot{p}_6 = \frac{b}{a} \frac{I_6}{I_2} \dot{p}_2$$

Substitute (8) into (7)

$$\dot{p}_2 = e_1(t) - \left(\frac{b}{a}\right)^2 \frac{R_5}{I_2} p_2 - \left(\frac{b}{a}\right)^2 \frac{I_6}{I_2} \dot{p}_2 - \frac{b}{a} \frac{1}{C_7} q_7 \tag{9}$$

(8)

 \square Solve (9) for \dot{p}_2

$$\dot{p}_2\left(\frac{I_2 + (b/a)^2 I_6}{I_2}\right) = e_1(t) - \left(\frac{b}{a}\right)^2 \frac{R_5}{I_2} p_2 - \frac{b}{a} \frac{1}{C_7} q_7 \tag{10}$$



Rearranging (10) gives the *first of two state equations*:

$$\dot{p}_{2} = -\frac{\left(\frac{b}{a}\right)^{2}R_{5}}{I_{2} + \left(\frac{b}{a}\right)^{2}I_{6}}p_{2} - \frac{\left(\frac{b}{a}\right)I_{2}}{\left(I_{2} + \left(\frac{b}{a}\right)^{2}I_{6}\right)C_{7}}q_{7} + \frac{I_{2}}{I_{2} + \left(\frac{b}{a}\right)^{2}I_{6}}e_{1}(t)$$
(11)

 $\square \text{ Next, move on to } \dot{q}_7$ $\dot{q}_7 = f_7 = f_4 = \frac{b}{a}f_3$ (12) $S_e \xrightarrow{e_1(t)} 1 \xrightarrow{3} TF \xrightarrow{4} 1 \xrightarrow{1}_{f_4 = b/a \cdot f_3} 1 \xrightarrow{\frac{1}{l_6} \cdot p_6} 1 \xrightarrow{\frac{1}{l_6} \cdot p_6} 1$ $I \xrightarrow{p_2} 1 \xrightarrow{f_4 = b/a \cdot f_3} 1 \xrightarrow{\frac{1}{l_7} \cdot q_7} 1 \xrightarrow{\frac{1}{p_6} \cdot q_8} 1 \xrightarrow{\frac{1}{p_6} \cdot$

The second state equation:

$$\dot{q}_7 = \frac{b}{a} \frac{1}{I_2} p_2$$

(9)

□ The state-space system model:

$$\begin{bmatrix} \dot{p}_2 \\ \dot{q}_7 \end{bmatrix} = \begin{bmatrix} -\frac{\left(\frac{b}{a}\right)^2 R_5}{I_2 + \left(\frac{b}{a}\right)^2 I_6} & -\frac{\left(\frac{b}{a}\right) I_2}{\left(I_2 + \left(\frac{b}{a}\right)^2 I_6\right) C_7} \\ \frac{b}{a} \frac{1}{I_2} & 0 \end{bmatrix} \begin{bmatrix} p_2 \\ q_7 \end{bmatrix} + \begin{bmatrix} \frac{I_2}{I_2 + \left(\frac{b}{a}\right)^2 I_6} \\ 0 \end{bmatrix} e_1(t)$$

□ With physical parameters:

$$\begin{bmatrix} \dot{p}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} -\frac{\left(\frac{b}{a}\right)^2 b}{m_1 + \left(\frac{b}{a}\right)^2 m_2} & -\frac{\left(\frac{b}{a}\right) m_1 k}{\left(m_1 + \left(\frac{b}{a}\right)^2 m_2\right)} \\ \frac{b}{a} \frac{1}{m_1} & 0 \end{bmatrix} \begin{bmatrix} p_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} \frac{m_1}{m_1 + \left(\frac{b}{a}\right)^2 m_2} \\ 0 \end{bmatrix} F_{in}(t)$$

Derivative causality in this case resulted from a modeling decision

- The lever was assumed to be rigid
- Adding some compliance to the lever arm eliminates derivative causality (see Section 3 notes)
 - Increases system model to *fourth-order*
 - Equation derivation simplified at the cost of model complexity
- In general, derive an expression for the energy variable of each energy-storage element in derivative causality
 - Multiple elements in derivative-causality will require solution of a system of equations