Stochastic cascade, symmetry and comparison method for the Navier-Stokes equations

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$$(\text{NSE}): \left\{ \begin{array}{ll} \partial_t u - \Delta u + u \cdot \nabla u + \nabla p = 0 & \text{ in } \mathbb{R}^d \times (0, \infty), \\ & \text{ div } u = 0 & \text{ in } \mathbb{R}^d \times (0, \infty), \\ & u(\cdot, 0) = u_0 & \text{ in } \mathbb{R}^d. \end{array} \right.$$

Goals:

- Brief overview on the global regularity and blowup.
- Le Jan-Sznitman's method for NSE on the Fourier side. Two problems: *stochastic explosion* and *integrability issue*.
- Conservation of frequency and applications to global regularity and finite time blowup of a model equation.

- Mild solutions are solutions obtained by Picard's iteration
 u_{n+1} = e^{Δt}u₀ + B(u_n, u_n). Regularity of u can be as good as U, e.g.
 smooth, analytic, etc.
- Local existence if u_0 is in scale-subcritical spaces. Global existence if u_0 is in scale-critical spaces.
- Finite time blowup is unknown. Finite time blowup, if exists, happens simultaneously in many spaces: $L_t^p L_x^q$, homogeneous Sobolev, Besov, etc.
- The energy identity plays the role of a mechanism to prevent blowup.
- Model equations that exhibit blowup: Montgomery-Smith (2001), Gallagher-Paicu (2009), Sinai-Li (2008), Tao (2015).

Fourier-transformed Navier-Stokes equations (FNS)

$$\hat{u}(\xi,t) = e^{-|\xi|^2 t} \hat{u}_0(\xi) + c_0 \int_0^t e^{-|\xi|^2 s} |\xi| \int_{\mathbb{R}^d} \hat{u}(\eta,t-s) \odot_{\xi} \hat{u}(\xi-\eta,t-s) d\eta ds$$

where $a \odot_{\xi} b = -i(e_{\xi} \cdot b)(\pi_{\xi^{\perp}}a)$.



Global regularity and blowup

- Li-Ozawa-Wang (2018): if supp û₀ ⊂ {(ξ₁,...,ξ_d) : ξ_k ≥ ℓ ≫ 1 ∀k} and ||û₀||_{L∞} ≤ ℓ then the solution is global-in-time.
- Feichtinger-Gröchenig-Li-Wang (2021): if supp û₀ ⊂ {(ξ₁,...,ξ_d) : ξ_k ≥ 0 ∀ k} and u₀ ∈ L² then the solution is global-in-time.

Theorem (P., 2021)

If supp $\hat{u}_0 \subset \mathbb{R}^d_+$ and \hat{u}_0 belongs to a Herz space $\dot{K}^{\alpha}_{p,q}$, where α, p, q are in a suitable range, then the solution is global-in-time. [If d = 2 or 3, L^2 is included in this family.]

$$\dot{K}^{\alpha}_{p,q} = \{f : |\xi|^{\alpha} f \mathbb{1}_{A_k} \in I^q L^p\}$$

Normalization: (LJS 1997, Bhattacharya et al 2003)

$$\begin{split} \chi(\xi,t) &= e^{-t|\xi|^2}\chi_0(\xi) \\ &+ \int_0^t e^{-s|\xi|^2}|\xi|^2\int_{\mathbb{R}^d}\chi(\eta,t-s)\odot_\xi\chi(\xi-\eta,t-s)H(\eta|\xi)d\eta ds \end{split}$$

where
$$\chi=c_0rac{\hat{u}}{h}$$
 and $H(\eta|\xi)=rac{h(\eta)h(\xi-\eta)}{|\xi|h(\xi)}.$

h satisfies $h * h = |\xi|h$, called majorizing kernel.

$$\chi = \mathbb{E}[\mathbf{X}_{FNS}]$$

where the stochastic process \boldsymbol{X}_{FNS} is defined recursively by. . .

Le Jan-Sznitman's construction of solutions

Consider the following event:



On this event,

$$\mathbf{X}_{\mathsf{FNS}}(\xi, t) = (\chi_0(W_{11}) \odot_{W_1} \chi_0(W_{12})) \odot_{\xi} \chi_0(W_2).$$

Three ingredients: clocks, branching distribution, product.

For $\chi = \mathbb{E}[\mathbf{X}_{FNS}]$ to be well-defined, one needs:

- *Stochastic non-explosion*: after finitely many branchings on the tree, all branches must terminate.
- Integrability: $\mathbb{E}|\mathbf{X}_{FNS}| < \infty$.

Can stochastic explosion happen?

•
$$h(\xi) = C_d |\xi|^{1-d}$$
: yes for $d = 3$, no for $d \ge 12$.

Overcome explosion issue: χ

$$\chi = \mathbb{E}[\mathbf{X}_{FNS} \mathbb{1}_{no.exp}]$$

Simplifying the product

How to deal with the integrability issue?

$$\begin{split} \chi(\xi,t) &= e^{-t|\xi|^2} \chi_0(\xi) \\ &+ \int_0^t e^{-s|\xi|^2} |\xi|^2 \int_{\mathbb{R}^d} \chi(\eta,t-s) \chi(\xi-\eta,t-s) H(\eta|\xi) d\eta ds \\ &= \mathbb{E}[\mathbf{X}(\xi,t)] \end{split}$$

where

$$\mathbf{X}(\xi,t) = \begin{cases} \chi_0(\xi) & \text{if } T_0 > t, \\ \mathbf{X}^{(1)}(W_1, t - T_0) \mathbf{X}^{(2)}(\xi - W_1, t - T_0) & \text{if } T_0 \leq t. \end{cases}$$

(MS):
$$\begin{cases} \partial_t u - \Delta u = \sqrt{-\Delta}(u^2) & \text{in } \mathbb{R}^d \times (0, \infty), \\ u(\cdot, 0) = u_0 & \text{in } \mathbb{R}^d \end{cases}$$

 χ_0 is a vectorial value function, yielding \mathbf{X}_{FNS} . $|\chi_0|$ is a scalar function, yielding \mathbf{X} .

• *Majorizing principle*:

$$|\mathbf{X}_{FNS}| \leq \mathbf{X}. \ (\mathsf{Thus}, \ \mathbb{E}|\mathbf{X}_{FNS}| \leq \mathbb{E}\mathbf{X}.)$$

• Power symmetry:

$$|\chi_0| \to |\chi_0|^{lpha}, \ \mathbf{X} \to \mathbf{X}^{lpha}.$$

• Conservation of frequency:

$$\chi_0 \to e^{\xi \cdot a} \chi_0, \quad \mathbf{X}_{FNS} \to e^{\xi \cdot a} \mathbf{X}_{FNS}, \quad \mathbf{X} \to e^{\xi \cdot a} \mathbf{X}.$$

Generalized majorizing principle:

• For a positive submultiplicative function ϕ ,

$$|\chi_0| \rightarrow \zeta_0 = \phi(|\chi_0|), \ \phi(\mathbf{X}) \le \mathbf{Z}$$

 $\bullet\,$ If, in addition, ϕ is convex then

$$|\chi_0| \rightarrow \zeta_0 = \phi(|\chi_0|), \ \phi(\chi) \le \zeta$$

Here **Z** and ζ are the stochastic process and solution yielded from ζ_0 .

Rule of thumb: the (Fourier) Montgomery-Smith equation controls the pointwise values of the solution to (Fourier) Navier-Stokes equations.

Frequency trap

If \hat{u}_0 is supported in a cone $C_{a,\theta} = \{\xi : \xi \cdot a \ge |\xi| |a| \cos \theta\}$ with an opening angle $\theta \in [0, \pi/2]$ then \hat{u} is supported in the same cone for all t.

Proof. $\chi_0 = \hat{u}/h$ is supported in the cone. For $\xi \in \mathbb{R}^d$ outside of the cone and with a realization of the cascade,

$$\mathbf{X} = \chi_0(w_1) \dots \chi_0(w_m).$$

Since $w_1 + \ldots + w_m = \xi$, at least one w_k has to lie outside of the cone. Thus, $\mathbf{X} = 0$.

Applications of the conservation of frequency

Conservation of frequency can be seen from the equation (Fourier side):

$$\hat{u}_0
ightarrow e^{\xi \cdot a} \hat{u}_0, \ \ \hat{u}
ightarrow e^{\xi \cdot a} \hat{u}$$

Global solution for supercritical/critical initial data

If supp $\hat{u}_0 \subset \mathbb{R}^d_+$ and $\hat{u}_0 \in L^2$, where d = 2 or 3, then the solution is global-in-time.

Proof. By Hölder's inequality,

$$\|e^{-\xi_d}\hat{u}_0\|_{\dot{K}_{1,2}^{-1}} \le \|\hat{u}_0\|_{\dot{K}_{2,2}^0}\|e^{-\xi_d}\mathbb{1}_{\xi_d\ge 0}\|_{\dot{K}_{2,\infty}^{-1}} <\infty.$$

Thus, $e^{-n\xi_d}\hat{u}_0$ is small in $\dot{K}_{1,2}^{-1}$ if *n* is large enough. $\dot{K}_{1,2}^{-1}$ is a scale-critical space for which global well-posedness is known (Cannone-Wu, 2012). Scale back to get the solution.

Global solution for subcritical initial data

If supp $\hat{u}_0 \subset \mathbb{R}^d_+$ and $\hat{u}_0 \in L^p$, where $1 \leq p < \frac{d}{d-1}$, then the solution is global-in-time.

Proof. The solution exists at least until time $T \sim \|\hat{u}_0\|_p^{-\delta}$ where $\delta = 2p/(d - (d - 1)p)$. The scaled initial data $\hat{u}_{0n} = e^{-n\xi_d}\hat{u}_0$ gives a solution at least until time $T_n \sim \|\hat{u}_{0n}\|_p^{-\delta} \to \infty$ as $n \to \infty$. Scale back to get a global solution. Blowup: solution $v \ge 0$ to the (Fourier) Montgomery-Smith equation is said to blow up after finite time if there exists T > 0 such that the set $\{\xi : v(\xi, T) = \infty\}$ has positive Lebesgue measure.

"Good" domains: a set \overline{D} (closure of an open set $D \subset \mathbb{R}^d$) is said to have Property (P) if any bounded function, compactly supported on \overline{D} , gives a global solution.

Theorem (P., 2021)

A half-space is a maximal set that has Property (P).

Proof. An open set strictly larger than a half-space must contain two balls symmetric w.r.t. the origin.

The half-space is optimal



Put $\Omega = B \cup (-B)$. *Key observation*: $|(\xi - \Omega) \cap \Omega| > c > 0$ for every $\xi \in A$. Put $\phi = \mathbb{1}_{\Omega}$ and choose the initial condition $v_0 = M\phi$. Oseen's representation of the solution says that

$$v = \sum_{k=1}^{\infty} M^k v_k$$
 where $v_1(\xi, t) = e^{-t|\xi|^2} \phi$,

The half-space is optimal

$$v_n(\xi, t) = \int_0^t \int_{\mathbb{R}^d} |\xi| e^{-(t-s)|\xi|^2} \sum_{k=1}^n v_k(\eta, s) v_{n-k}(\xi - \eta, s) d\eta ds$$

Put $w_k(t) = \inf_{\xi \in A} v_{2^k}(\xi, t)$. By induction,

$$w_k(t)\gtrsim e^{-lpha 2^k}2^k(te^{-t})^{2^k-1}$$
 $\forall k$

For $\xi \in A$,

$$\mathsf{v}(\xi,1) \geq \sum M^{2^k} \mathsf{v}_{2^k}(\xi,1) \gtrsim \sum M^{2^k} e^{-lpha 2^k} 2^k e^{-(2^k-1)} o \infty.$$

Probability for all paths to end up in the half-space

Branching distribution:



 $p_k(\xi, t) = \mathbb{P}_{\xi}(by \text{ time } t, \text{ all paths ending up in } \mathbb{R}^d_+)$

Take $\chi_0(\xi) = M \mathbb{1}_{\xi_d \ge 0}$, or equivalently $\hat{u}_0 \sim M |\xi|^{1-d} \mathbb{1}_{\xi_d \ge 0}$. Note that $e^{-\xi_d} \hat{u}_0 \in L^p$ for any 1 . $Global solution <math>\hat{u} \sim |\xi|^{1-d} \chi$ exists. Thus,

$$\chi = \sum M^k p_k < \infty \quad \forall M > 0.$$

 p_k must decay faster than any exponential function of k. To quantify the rate of decay, one needs to understand more the geometry of the half-space.

$$p_n(\xi, t) = \int_0^t \int_{\mathbb{R}^d} |\xi|^2 e^{-(t-s)|\xi|^2} \sum_{k=1}^n p_k(\eta, s) p_{n-k}(\xi - \eta, s) H(\eta|\xi) d\eta ds$$

Probability for all paths to end up in the half-space

The key idea is the following:

$$\int_{(\xi-\mathbb{R}^d_+)\cap\mathbb{R}^d_+} H(\eta|\xi) \leq c_{\delta,d} \left(rac{\xi_d}{|\xi|}
ight)^{\delta} \quad orall \delta \in (0,1) \, .$$



The following lemma is needed for the induction procedure: for $\phi(k) = (k!)^{\gamma}$, where $\gamma \in (0, 1)$, one has

$$\sum_{k=1}^{n-1} \frac{a^k}{\phi(k)} \frac{b^{n-k}}{\phi(n-k)} \leq \frac{(a+b)^n}{\phi(n)} n^{\gamma-1} \quad \forall a, b > 0, \ n \in \mathbb{N}.$$

Theorem (P., 2021)

$$p_k(\xi,t) \leq c_d^{k-1} rac{(\xi_d \sqrt{t})^{rac{1}{2}(k-1)}}{\{(k-1)!\}^{1/8}} e^{-|\xi|\sqrt{t}}.$$

Thank You!